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# HW #8

## 1. Coupled Harmonic Oscillators

The Hamiltonian is

$$H = \frac{p_1^2}{2M} + \frac{p_2^2}{2M} + \frac{p_3^2}{2m} + \frac{1}{2}k(x_3 - x_1 - d)^2 + \frac{1}{2}k(x_2 - x_3 - d)^2$$
For the sake of simplicity, we redefine  $x_1 \to x_1 - d$ ,  $x_2 \to x_2 + d$  and the Hamiltonian is 
$$H = \frac{p_1^2}{2M} + \frac{p_2^2}{2M} + \frac{p_3^2}{2m} + \frac{1}{2}k(x_3 - x_1)^2 + \frac{1}{2}k(x_2 - x_3)^2$$

#### (a) Born-Oppenheimer Approximation

Assume 
$$x_1$$
 and  $x_2$  are fixed. Then the potential energy for  $x_3$  is
$$\frac{1}{2}k(x_3^2 - 2x_3x_1 + x_1^2 + x_2^2 - 2x_2x_3 + x_3^2) = \frac{1}{2}k(2x_3^2 - 2(x_1 + x_2)x_3 + x_1^2 + x_2^2)$$

$$= \frac{1}{2}2k(x_3 - \frac{x_1 + x_2}{2})^2 - k\frac{(x_1 + x_2)^2}{4} + \frac{1}{2}k(x_1^2 + x_2^2)$$

$$= \frac{1}{2}2k(x_3 - \frac{x_1 + x_2}{2})^2 + \frac{1}{4}k(x_1 - x_2)^2$$

The energy levels for the  $x_3$  degree of freedom is therefore  $(n + \frac{1}{2}) \hbar \sqrt{\frac{2k}{m}}$ .

Taking the ground state for the  $x_3$  degree of freedom, the Hamiltonian reduces to

$$H_{\text{eff}} = \frac{p_1^2}{2M} + \frac{p_2^2}{2M} + \frac{1}{2} \hbar \sqrt{\frac{2k}{m}} + \frac{1}{4} k(x_1 - x_2)^2$$

 $H_{\text{eff}} = \frac{p_1^2}{2M} + \frac{p_2^2}{2M} + \frac{1}{2} \hbar \sqrt{\frac{2k}{m}} + \frac{1}{4} k(x_1 - x_2)^2.$ Separating the center-of-mass motion with  $P = p_1 + p_2$ ,  $X = \frac{x_1 + x_2}{2}$ ,  $p = \frac{p_1 - p_2}{2}$ ,  $x = x_1 - x_2$ , it becomes

$$H_{\rm eff} = \frac{P^2}{4M} + \frac{p^2}{M} + \frac{1}{4} k x^2 + \frac{1}{2} \hbar \sqrt{\frac{2k}{m}}$$
. Leaving the trivial translational energy aside, the energy levels are therefore

$$E = (n' + \frac{1}{2}) \hbar \sqrt{\frac{k}{M}} + \frac{1}{2} \hbar \sqrt{\frac{2k}{m}}.$$

## (b) Exact Solution

We identify the normal modes of the coupled harmonic oscillators.

$$L = \frac{1}{2} M \dot{x}_1^2 + \frac{1}{2} M \dot{x}_2^2 + \frac{1}{2} m \dot{x}_3^2 - \frac{1}{2} k(x_3 - x_1)^2 - \frac{1}{2} k(x_2 - x_3)^2.$$

We rescale the coordinates to make the kinetic terms the same for all three,  $x_1 = \frac{1}{\sqrt{M}} q_1$ ,  $x_2 = \frac{1}{\sqrt{M}} q_2$ ,  $x_3 = \frac{1}{\sqrt{m}} q_3$ . Then

$$L = \frac{1}{2} \dot{q_1}^2 + \frac{1}{2} \dot{q_2}^2 + \frac{1}{2} \dot{q_3}^2 - \frac{1}{2} k \left( \frac{1}{\sqrt{m}} q_3 - \frac{1}{\sqrt{M}} q_1 \right)^2 - \frac{1}{2} k \left( \frac{1}{\sqrt{m}} q_3 - \frac{1}{\sqrt{M}} q_2 \right)^2.$$

We can write the potential term as

$$V = \frac{1}{2} k(q_1, q_3, q_2) \begin{pmatrix} 1/M & -1/\sqrt{mM} & 0 \\ -1/\sqrt{mM} & 2/m & -1/\sqrt{mM} \\ 0 & -1/\sqrt{mM} & 1/M \end{pmatrix} \begin{pmatrix} q_1 \\ q_3 \\ q_2 \end{pmatrix}.$$

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$$In[5] := \text{ Eigenvalues} \left[ \left\{ \left\{ 1 \, / \, M \, , \, -1 \, \middle/ \, \sqrt{m \, M} \, , \, 0 \right\}, \, \left\{ -1 \, \middle/ \, \sqrt{m \, M} \, , \, 2 \, / \, m \, , \, -1 \, \middle/ \, \sqrt{m \, M} \right\}, \, \left\{ 0 \, , \, -1 \, \middle/ \, \sqrt{m \, M} \, , \, 1 \, / \, M \right\} \right\} \right]$$

$$Out[5] = \left\{ 0 \, , \, \frac{1}{M} \, , \, \frac{m + 2 \, M}{m \, M} \right\}$$

Therefore, it is given by three independent harmonic oscillators of frequencies 0,  $\sqrt{k/M}$ , and  $\sqrt{k(2M+m)/mM}$ . The zero-frequency mode corresponds to the overall translation, while the other two give the harmonic oscillator levels,  $E = \left(n' + \tfrac{1}{2}\right)\hbar\sqrt{\tfrac{k}{M}} + \left(n + \tfrac{1}{2}\right)\hbar\sqrt{\tfrac{k(2M+m)}{mM}} \,.$ 

#### (c) Comparison

Taking n = 0 in the exact result, we find

$$E = (n' + \frac{1}{2}) \hbar \sqrt{\frac{k}{M}} + \frac{1}{2} \hbar \sqrt{\frac{k(2M+m)}{mM}},$$
 while the Born-Oppenheimer approximation gives

$$E = (n' + \frac{1}{2}) \hbar \sqrt{\frac{k}{M}} + \frac{1}{2} \hbar \sqrt{\frac{2k}{m}}.$$

Therefore the Born-Oppenheimer approximation gives the correct result up to a correction suppressed by m/M as expected.

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### 2. A=14

Looking at the nitrogen, the levels with spin/parity that appear together in carbon and nitrogen with similar excitation energies are identified as I = 1, wille those that appear only in nitrogen as I = 0.

From the ground state and up,

```
J^P
       I
1+
       0
0^{+}
       1
1+
       0
0-
       0
2-
       0
1-
       0
3-
       0
1+
       0
3+
       0
2+
       0
2-
       0
               8062 - 2313 = 5749, in rough accordance with the excitation energies in carbon and oxygen
1-
       1
4-
       0
0^{+}
       1
0-
       0
3-
       1
5-
       0
2+
       0
               appears too close to 3^- with I = 1 compared to carbon and oxygen
3+
       0
2+
               9172 - 8907 = 265, in rough accordance with the 2^+ - 3^- splittings in carbon and oxygen
       1
etc
```